

Performance Comparison of Long-Wave Infrared Imaging Spectrometer Between Dyson Form and Offner Form

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Abstract: In view of the difficulties in traditional long-wave infrared imaging spectrometer which is hard to realize a high signal-to-noise ratio and miniaturization as well under the weak remote sensing signal, Offner convex grating spectrometer and Dyson concave grating spectrometer, both having concentric structure, are designed and analyzed in the band of 8-12 μm . The diffraction angle expressions of the two spectrometers are obtained and the diffraction characteristics are acquired. Both of the spectrometers are designed in Zemax environment under different F-numbers and different grating constants with the same slit, spatial resolution, spectral resolution and detector. The results show that Dyson grating spectrometer possesses the advantages of higher throughput and smaller volume, and Offner grating spectrometer possesses the advantage of more accessible material and the absence of chromatic aberration. The differences between Dyson form and Offner form show that the former is a better choice in the long-wave infrared imaging spectrometer.

Keywords: long-wave infrared; imaging spectrometer; grating constant; concentric configuration

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1 Introduction

Offner form^[1-2] and Dyson form^[3-4] are the two typical concentric configurations. Grating spectrometers based on concentric configuration can offer appreciable advantages over traditional arrangements, such as high numerical aperture, simple and compact structure, and low geometric distortions^[5-7]. In visible and near infrared system, the Offner form has received more attention than the Dyson form. The VIRTIS (visible infrared thermal imaging spectrometer) and the M³ (moon mineralogy mapper) imaging spectrometer are all based on Offner form^[8-10]. Dyson form is usually introduced into the system which needs higher throughput such as the PRISM (portable remote imaging spectrometer)^[11]. The band of long-wave infrared (LWIR, nominally 8-12 μm) is one of the atmosphere windows, and the spectral information from this wavelength range is extremely valuable for earth science research. The Mako airborne sensor with a fast F-number of 1.25 is designed based on Dyson form works in this band and it can realize spatial resolution of 0.55 mrad and spectral resolution of 44 nm^[12-13].

Most of the domestic studies about concentric

imaging spectrometers are in the laboratory stage^[14]. A portable imaging spectrometer based on Offner form was developed in 2010. It worked in the band of 0.4 – 0.8 μm with the F-number of 3.5 and the spectral resolution of 2.1 nm^[15]. An imaging spectrometer based on Dyson was designed in 2014. It worked in the band of 7.5–10 μm with the F-number of 1.2 and the spectral resolution of 50 nm^[16].

We have studied the grating spectrometers based on Offner form and Dyson form. Then the difference in terms of F-number is discussed and their diffraction angle expressions are given. Considering that the throughput and the system volume are important in LWIR spectrometer, we choose them as indexes to evaluate the performance. Firstly, the comparison is made under the condition of the same focal length but different F-number. Then, the comparison is made under the condition of same F-number but different focal length. The specific designs about the Offner form and Dyson form are presented, and the advantages and disadvantages of them are summarized.

2 General Theory in Grating Spectrometer Design

The optical system of grating imaging spectrometer

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can be divided into two main subsystems including telescope objective and grating spectrometer. The second part, which is the core part, will be discussed in details.

Grating spectrometer obeys the diffraction function

$$d(\sin i \pm \sin \theta) = \pm m\lambda \quad (1)$$

where d is the grating constant; i is the incident angle; θ is the diffraction angle; m is the diffraction order and λ is the wavelength. The dispersion angle can be expressed as

$$d\theta = \frac{m}{d\cos\theta}d\lambda \quad (2)$$

When the focal length of focusing lens equals to f' , the dispersion can be expressed as

$$\frac{dl}{d\lambda} = \frac{mf'}{d\cos\theta} \quad (3)$$

The dispersion maintains linear within a relative narrow band of wavelengths which benefits from the stable diffraction angle. It is shown in Eqs. (1) and (2) that the diffraction angle can be larger at longer wavelength. Steep angle leads to relative higher dispersion, and relative high dispersion can bring higher off-axis aberration. Diffraction angle can be smaller for a relative larger d , and so the off-axis aberration can be lower. It can be seen from Eq. (3) that maintaining $dl/d\lambda$ invariant, the focal length increases as d increases which enlarges the system volume. The analyses above demonstrate that grating constant has effects on imaging quality and system volume, and the imaging quality and the system volume are mutually restricted.

In remote sensing system, spatial resolution (Δl) can be expressed as

$$\Delta l = \frac{l_{\text{slit}}}{f_{\text{tele}}}H$$

where l_{slit} is the width of slit; f_{tele} is the focal length of the telescope objective and H is the height of the flying platform (aircraft or satellite). Δl and H are usually required, and so a thin l_{slit} is needed to guarantee small volume. But too thin l_{slit} will decrease the SNR (signal to noise ratio). When the width of monochromatic image of slit (l'_{slit}) equals the size of pixel (a), the system has the best spectral resolution. Taking the volume, SNR and spectral resolution into consideration, we can get a reasonable design when

$$l_{\text{slit}} = l'_{\text{slit}} = a \quad (4)$$

Spectral resolution of the detector ($\delta\lambda$) is

$$\delta\lambda = \frac{\lambda_2 - \lambda_1}{N} \quad (5)$$

where N is the number of spectral dimension pixels, and $\lambda_2 - \lambda_1$ is the range of wavelength. According to Eq. (3), spectral resolution determined by the spectrometer can be expressed as

$$\Delta\lambda = \frac{d\cos\theta}{mf'}dl$$

In an ideal imaging system, $\delta\lambda$ is equal to $\Delta\lambda$, and dl is equal to a .

In order to guarantee the system having enough energy responsibility, we choose cooled HgCdTe with pixel array of 320×256 and pixel size of $30 \mu\text{m}$ as the detector. So the width of slit should be equal to $30 \mu\text{m}$ according to Eq. (4). Then the slit size is $9.6 \text{ mm} \times 30 \mu\text{m}$, width of spectrum is 7.68 mm , and the system spectral resolution, according to Eq. (5), will be 15.6 nm .

3 Introductions of the Two Concentric Grating Spectrometers

3.1 Offner Convex Grating Spectrometer

Offner form consists of a concave mirror (M_1) and a convex mirror (M_2) with radius of R_1 and R_2 respectively, and the two mirrors have a common center C as shown in Fig. 1, where u is the aperture angle and h is the height of the object. M_2 is the system stop and is located around the focal plane of M_1 . Therefore, the design is telecentric. It is obvious that Offner form is symmetrical and its magnification is -1 . Spherical aberration and astigmatism are the only two aberrations needing to be considered, benefiting from the symmetry. Besides, there is no need to worry about chromatic aberration, because all optical elements are mirrors. Only when the concentricity and the equation $R_2 = R_1/2$ are fulfilled simultaneously^[17-19], the system have the least aberration. But Offner form is still limited by spherical aberration, and the inclination of the meridional image field at h is approximately proportional to $1/F^3$.

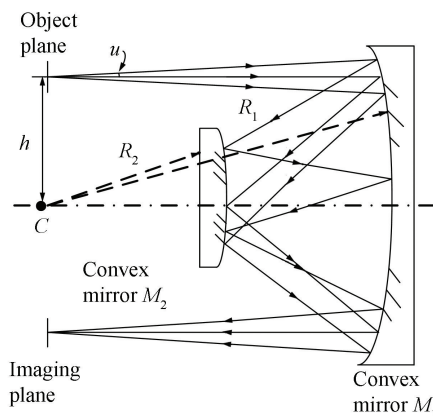


Fig. 1 Schematic of Offner concentric system

Replacing M_2 by a convex reflection grating (G) will get Offner convex grating spectrometer. The diffraction characteristics of it can be obtained according to Fig. 2, where l is the length from point C to point D , and α is the incident angle of M , and β is

the angle forming by optical axis and reflective ray from M , and i_1 and θ_1 is the incident angle and diffraction angle of G respectively.

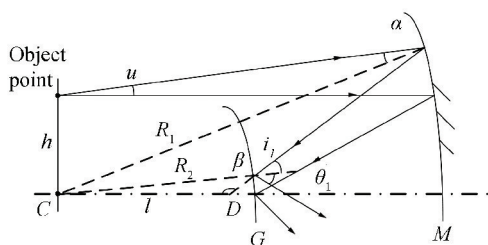


Fig. 2 Diffraction schematic of Offner spectrometer

The following equations are acquired according to the sine theorem referred to Fig. 3.

$$\frac{\sin \alpha}{h} = \frac{\sin(90^\circ + u)}{R_1}$$

$$\frac{\sin \beta}{R_1} = \frac{\sin \alpha}{l}$$

$$\frac{\sin i_1}{l} = \frac{\sin \beta}{R_2}$$

and hence

$$\sin \alpha = \frac{\cos u}{R_1} h \quad (6)$$

$$\sin i_1 = \frac{\sin \alpha}{R_2} R_1 \quad (7)$$

In order to avoid obscuration of rays, h should be

$$h = R_1 \tan u \quad (8)$$

Offner spectrometer obeys the function

$$d(\sin i_1 - \sin \theta_1) = -\lambda \quad (9)$$

From Eqs. (6)–(9), θ_1 can be acquired:

$$\theta_1 = \arcsin\left(\frac{\lambda}{d} + \frac{R_1 \sin u}{R_2}\right) \quad (10)$$

3.2 Dyson Concave Grating Spectrometer

Dyson form consists of a plano-convex lens and a concave mirror, and the lens and the mirror share a common center C , as shown in Fig. 3. The concave mirror is the system stop and it is located around the focal plane of the plano-convex lens. The system is a telecentric system with magnification of +1. So there is a relation between R and r as

$$R = \frac{n}{n-1} r$$

where n is the refractive index.

Dyson form is free from all Seidel aberrations without restriction of aperture, but it also suffers from high order aberrations. The aberrations will cause a distance of s departing from the ideal position of the image along the axis, and s can be expressed as

$$s = \frac{n(n-1)h^4}{r^3} \cos^3 u$$

It is obvious that the u makes less contribution to $s^{[3]}$.

Dyson form is catadioptric, and we have to select applicable material to fabricate the plano-convex lens. The material should possess high transmission over 8–12 μm and should have high Abbe number to reduce the chromatic aberration. We select ZnSe because its Abbe number is higher (58.61), and its absorption coefficient is lower (<0.0005).

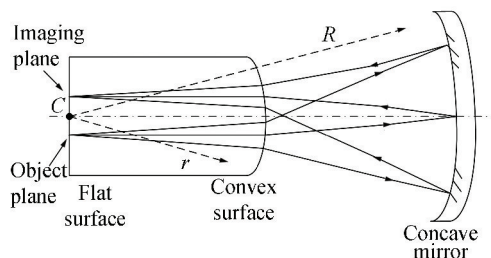


Fig. 3 Schematic of Dyson concentric system

It will be Dyson concave grating spectrometer after substituting the concave mirror with a concave reflection grating. The diffraction characteristics of Dyson spectrometer can be obtained according to Fig. 4, where θ_2 represents diffraction angle.

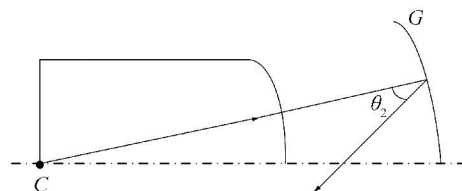


Fig. 4 Diffraction schematic of Dyson spectrometer

Dyson spectrometer obeys the following function

$$d(\sin i_2 + \sin \theta_2) = \lambda$$

Obviously incident angle i_2 equals to 0, the diffraction angle θ_2 is as

$$\theta_2 = \arcsin\left(\frac{\lambda}{d}\right) \quad (11)$$

From the descriptions above, we can infer the characteristics of Offner form and Dyson form as: (1) Offner form is not suitable in lower F-number system while Dyson form is, (2) In comparison of Eq. (10) to Eq. (11), it can be seen that θ_1 is bigger than θ_2 , and so the off-axis aberration of Offner is higher, and (3) According to Eq. (10), θ_1 is bigger in the system with large numerical aperture.

F-number has effect on the throughput and grating constant has effect on the volume, and the throughput and the volume are both important parameters in LWIR imaging spectrometer. Therefore, we compare Offner form and Dyson form in different F-numbers and different grating constants.

4 Comparison in Different F-numbers

The performance of Offner spectrometer is limited by aperture angle, and it is said that NA should be less

than 0.18 if we want to get good imaging ^[20]. The relationship between the numerical aperture of NA and the F-number of $F/\#$ is as

$$NA \approx \frac{1}{2F/\#} \tag{12}$$

$F_1/\#$ approximately equals 2.8 when $NA = 0.18$ according to Eq. (12). Because SNR is proportional to $1/(F/\#)^2$ theoretically in infrared system, we set $F_2/\#$ to be 2 which is $2^{-1/2}$ times of $F_1/\#$. The grating pitch can be actually manufactured quite small, but d must be higher than $18.75\text{ }\mu\text{m}$ to enable θ_1 less than 90° according to Eq. (11). At last, we pick up the grating constants of $100\text{ }\mu\text{m}$ and $40\text{ }\mu\text{m}$ in our designs in terms of the off-axis aberration in Offner form.

All of the parameters needed have been provided. Then we compare Offner form and Dyson form in different F-numbers firstly, and all of the indicators are listed in Table 1.

Table 1 Indicators for spectrometer design

Performance parameters	Values
Spectral range(μm)	8–12
Entrance slit size(mm \times μm)	9.6 \times 30
Width of spectrum(mm)	7.68
Spectral resolution(nm)	15.6
$F_1/\#(F_2/\#)$	2.8(2)
$d_1(\mu\text{m})$	100
$f_1'(\text{mm})$	192
Diffraction order in Offner (Dyson)	-1(+1)

For comparison, we design two Dyson spectrometers with spherical surfaces, as shown in Fig. 5. From Fig. 5, we can see that each wavelength can be focused well, which implying both of them can realize ideal imaging. Fig. 5(a) is with F-number of 2.8, and Fig. 5(b) is with F-number of 2. Both of the spectrometers can realize spectral resolution of 15.6 nm , and are with similar volume.

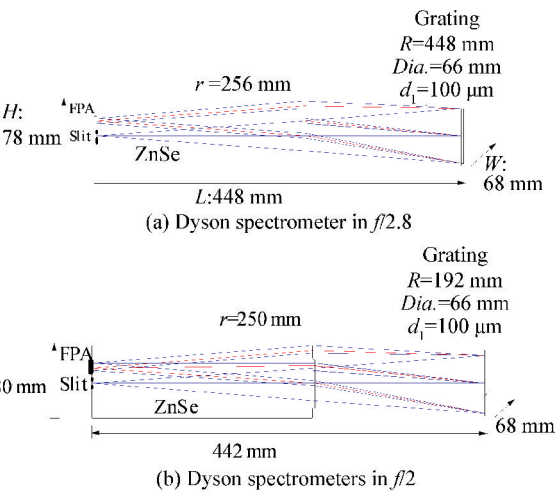


Fig. 5 Two Dyson spectrometers with different F-number

Then we design two Offner spectrometers. Fig. 6 shows the designed spectrometer with F-number of 2.8. It can be seen that the ideal imaging can be realized. With this spectrometer, spectral resolution of 15.6 nm can be realized. The one with F-number of 2 fails to perfect imaging due to large aberrations, especially spherical aberration. Fig. 7 shows geometric spot diagrams of the designed system, where the center circles show Airy disk. It can be seen that the spots are much bigger than Airy disks, implying unacceptable imaging quality.

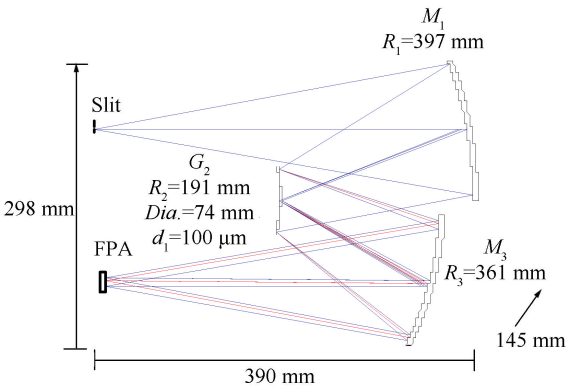


Fig. 6 Offner spectrometer in f/2.8

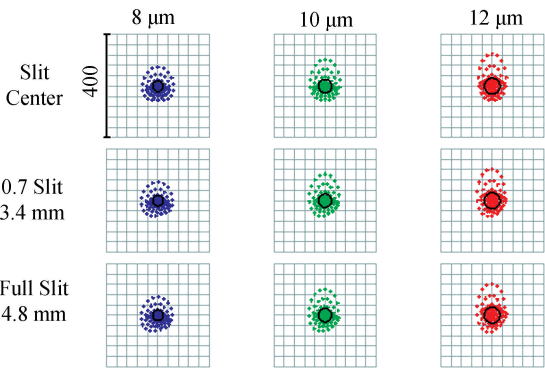


Fig. 7 Geometric spot diagrams for f/2 Offner spectrometer

We make suitable decenter of M_1 , G_2 and M_3 from the axis vertically, and then set M_1 and M_3 as aspherical surfaces. A better system is obtained, and Fig. 8 shows the geometric spot diagrams of the system. It can be seen that the sizes of spots are improved greatly, but still larger than Airy disks.

Obviously, Offner form is suitable only for the system with larger F-number, while Dyson form can meets the requirement of system with smaller F-number. The spherical aberration existing in Offner spectrometer with smaller F-number cannot be corrected perfectly. With the same F-number, Dyson form owns a smaller volume because of its catadioptric optical structure. On the other hand, Offner form is of reflective optical structure, inevitably resulting in longer off-axis length, and this makes its volume bigger.

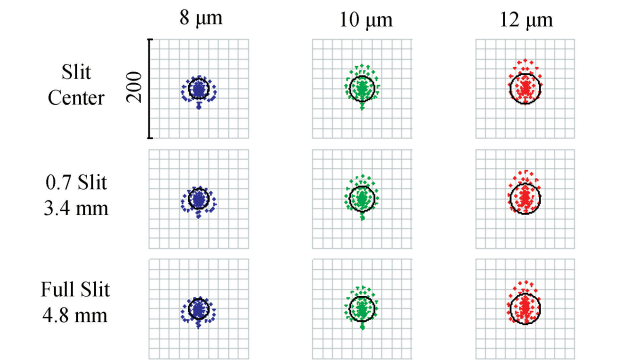


Fig. 8 Geometric spot diagrams for the corrected Offner spectrometer

5 Comparison in Different Grating Constants

In this part, we will compare Offner form and Dyson form in different grating constants, and all of the indicators are listed in Table 2.

Table 2 Indicators for spectrometer design	
Performance parameters	values
Spectral range(μm)	8–12
Entrance slit size(mm× μm)	9.6×30
Width of spectrum(mm)	7.68
Spectral resolution(nm)	15.6
$F_1/\#$	2.8
d_2 (μm)	40
f_2' (mm)	76.8
Diffraction order in Offner (Dyson)	-1(+1)

For comparison, we design two concentric spectrometers of a lower d with spherical surfaces. Dyson form can realize ideal imaging, as shown in Fig. 9. The spectral resolution can still reach 15.6 nm. Offner spectrometer can realize ideal imaging as well, but its spectral resolution cannot be 15.6 nm because of the large geometric spots as shown in Fig. 10.

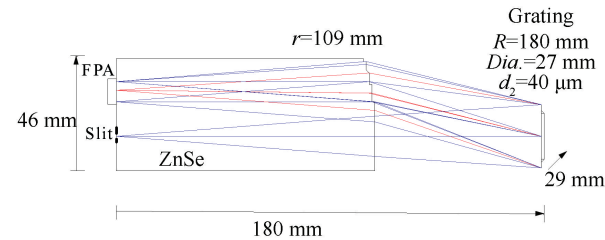


Fig. 9 Dyson spectrometer with d_2 of 40 μm

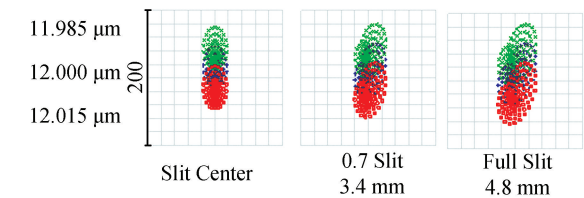


Fig. 10 Geometric spots of Offner spectrometer with d_2 of 40 μm

In order to correct the aberrations, we make suitable decenter of M_1 , G_2 and M_3 from the axis vertically, and then set M_1 and M_3 as aspherical surfaces as before. The designed structure is shown in Fig. 11. Its spectral resolution can reach 15.6 nm.

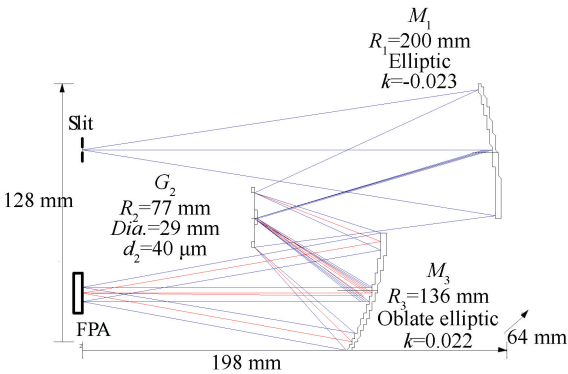


Fig. 11 Corrected Offner spectrometer with d_2 of 40 μm

Design results clearly show that, just with spherical mirror, Dyson spectrometer has better imaging quality with a lower grating constant, and therefore higher spectral resolution with smaller size of system can be realized. Although imaging quality of Offner form with lower grating constant is worse, it can be improved by decentering and aspherizing.

6 Design of Dyson Spectrometer with Better Performances

It can be inferred that Dyson form has ideal imaging in lower F-number and lower grating constant from the comparisons. In practice, the packaged focal plane in chip carriers should be laterally separated from the slit, and more importantly, the wire bonds and protective windows make it necessary that there should be an axial gap between the image plane and the rear face of Dyson lens [21–22]. In this case, Dyson form may lose its advantage of perfectly imaging. Spherical aberration is the main aberration and so we introduce an aspheric corrector nearby the stop. We design a Dyson spectrometer with $f/2$ and grating constant of 40 μm with corrector lens, as shown in Fig. 12.

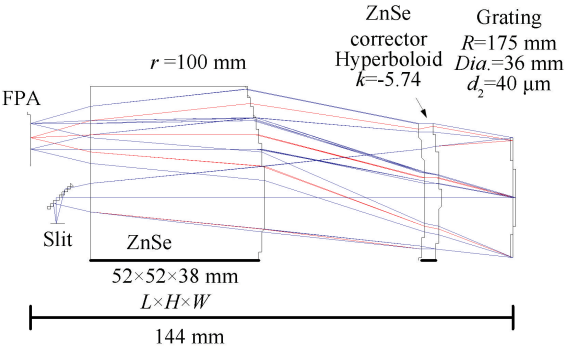


Fig. 12 Dyson spectrometer in $f/2$ and d of 40 μm with corrector lens

Finally, we have designed a cooled Dyson LWIR imaging spectrometer. We choose three-mirror off-axis aspherical optical system as the fore-optics and adopt re-imaging to realize 100% cold stop efficiency. The imaging spectrometer is with the size of 300 mm × 250 mm×116 mm, the spectral resolution of 25 nm and the spatial resolution of 0.2 mrad.

7 Conclusions

Two typical concentric spectrometers are

designed, and both of them possess the advantages of simple and compact structure. Compared to Offner form, Dyson spectrometer can realize ideal imaging in lower F-number with even smaller volume just with spherical mirror. Thus it can meet the requirements of high signal-to-noise ratio and miniaturization in LWIR imaging spectrometer. In terms of Offner spectrometer, the realization of miniaturization is easier as compared with the realization of high throughput. All of the comparisons are tabulated in Table 3.

Table 3 The comparison of Offner form and Dyson form

Form	Structure	Material	Chromatic aberration	F-number	Volume
Dyson	Simple and compact	Limited	Existent	Lower	Smaller
Offner	Simple and compact	Accessible	Absent	Higher	Bigger

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